

**Cavity magnonics** 

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- Intro: Spin waves and magnons
- Magnet-cavity coupling
- Level attraction
- Chiral magnon propagation
- Quantum magnonics

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# **Cavity magnonics**

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- Artem Bondarenko
- Enes Ilbuga
- YMB

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#### - Chuanpu Liu - Haiming Yu

- University of Manitoba
- Bimu Yao
- Can-Min Hu

- Iran University of Science and Technology
  - Babak Zare Rameshti

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### **Larmor precession**

#### Magnetic moment in external field

Torque:  $\mathbf{\tau} = \mathbf{M} \times \mathbf{B}$ 



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#### Magnons are elementary excitations of magnetic structure

#### Classical limit (large occupation numbers): spin waves



#### Image credit: Jens Böning, Wikimedia Commons

Spin Hamiltonian (simplest):  $\hat{H}_S = -J \sum_{\langle ij \rangle} \hat{S}_i \cdot \hat{S}_j - g \mu_B B \cdot \sum_i \hat{S}_i$ 

Spin wave spectrum of an isotropic 1D FM chain:  $\hbar\omega(k) = 2JS \left(1 - \cos ka\right) + g\mu_B B$ 

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# Imaging of spin waves

### Spin waves in YIG films imaged by NV center magnetometry



I. Bertelli, J. J. Carmiggelt, T. Yu, B. G. Simon, C. C. Pothoven, G. E. W. Bauer, YMB, J. Aarts, and T. van der Sar, Science Adv. **6**, eabd3556 (2020).

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# Yttrium Iron Garnet – ferrimagnetic insulator with the highest magnetic quality

Gilbert damping parameter:





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#### Very recently: also vanadium tetracyanoethylene

Gilbert damping comparable to YIG

Q. Hu ... G . Fuchs, arXiv:2212.04423





#### Surface spin waves in YIG (Damon-Eschbach modes) Technische Universiteit Delf

YIG: Dispersion relation is very anisotropic

Two competing mechanisms: Exchange and dipolar interactions

Spin waves travelling perpendicular to the magnetization







# **Quantization of spin waves**

#### Holstein-Primakoff transformation

$$egin{aligned} \hat{S}_{+} &= \hbar\sqrt{2S}\sqrt{1-rac{\hat{m}^{\dagger}\hat{m}}{2S}}\hat{m} \ \hat{S}_{-} &= \hbar\sqrt{2S}m^{\dagger}\sqrt{1-rac{\hat{m}^{\dagger}\hat{m}}{2S}} \ \hat{S}_{z} &= \hbar\left(S-\hat{m}^{\dagger}\hat{m}
ight) \end{aligned}$$

#### Plane waves:

$$\hat{m}_k = \frac{1}{\sqrt{N}} \sum_{R_i} \exp\left(-ik \cdot R_i\right) \hat{m}_k$$

#### Linearized transformation:



#### One-mode linear Hamiltonian:

$$\hat{H} = \hbar \omega_m \hat{m}^\dagger \hat{m}$$





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### **Microwave cavities**



Copper box cavity: A. Bienfat .... P. Bertet, Nature **531** 74 (2016)



Lumped-element LC resonator: F. Yoshihara .... K. Semba, Nature Physics **13** 44 (2017)



### **Microwave cavities**

#### Generally: Multimode cavities

Quantize vector potential (homogeneous situation):

$$\begin{split} \mathbf{A}(\mathbf{r},t) &= \mathbf{A}^{+}(\mathbf{r},t) + \mathbf{A}^{-}(\mathbf{r},t), \quad \mathbf{A}^{+} = (\mathbf{A}^{-})^{\dagger} \\ \mathbf{A}^{+}(\mathbf{r},t) &= \sum_{k} u_{k}(\mathbf{r}) \hat{a}_{k} e^{-i\omega_{k}t}, \omega_{k} = ck/n \quad \left(\nabla^{2} + k^{2}\right) u_{k} = 0 \\ \\ \text{Free cavity Hamiltonian:} \quad \hat{H} &= \sum_{k} \hbar \omega_{k} a_{k}^{\dagger} \hat{a}_{k} \\ \\ \text{Fields:} \quad \hat{\mathbf{E}}^{+}(\mathbf{r},t) &= i \sum_{k} \sqrt{\frac{\hbar \omega_{k}}{2V \varepsilon \varepsilon_{0}}} u_{k}(\mathbf{r}) \hat{a}_{k} e^{-i\omega_{k}t} \\ \\ \quad \hat{\mathbf{B}}^{+}(\mathbf{r},t) &= i \sum_{k} \sqrt{\frac{\hbar}{2V \varepsilon \varepsilon_{0} \omega_{k}}} \nabla \times u_{k}(\mathbf{r}) \hat{a}_{k} e^{-i\omega_{k}t} \\ \\ \\ \text{Mode volume:} \quad V_{k} &= \frac{\int |E_{k}|^{2} dV}{\max |E_{k}|^{2}} \end{split}$$

In most situations: Want to work with one mode

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# **Magnet-cavity Interaction**

Mechanism: interaction of magnetization with the cavity field  $M \cdot B$ Hamiltonian of an interaction of a single cavity mode with a single (almost resonant) magnon mode:

$$\hat{H} = \hbar\omega_c \hat{a}^{\dagger} \hat{a} + \hbar\omega_m \hat{m}^{\dagger} \hat{m} + g \left( \hat{a}^{\dagger} \hat{m} + \hat{a} \hat{m}^{\dagger} \right)$$



Strong coupling in a cavity predicted by Soykal and Flatte, Phys. Rev. Lett. **104**, 077202 (2010)

g – coupling constant Strong coupling regime means:  $\kappa_c, \kappa_m \ll g \ll \omega_c, \omega_m$ O Yaroslav M. Blanter





### Interaction

#### YIG film in a cavity

#### Hübl et al, Phys. Rev. Lett. **111**, 127003 (2013)



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### Interaction

#### Tabuchi et al, Phys. Rev. Lett. **113**, 083603 (2014)



Normal mode splitting between a magnon (YIG sphere) and a cavity mode

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## **Magnon spintronics**

Spin waves can carry information:

Ferromagnetic metals: Spin current is carried by electrons – Ohmic dissipation

Ferromagnetic insulators: Spin current carried by spin waves – Weak intrinsic damping of spin waves We want to be able to excite, manipulate, and read out spin waves

*Cavity Magnonics*, Babak Zare Rameshti, Silvia Viola Kusminskiy, James A. Haigh, Koji Usami, Dany Lachance-Quirion, Yasunobu Nakamura, Can-Ming Hu, Hong X. Tang, Gerrit E. W. Bauer, and YMB, Physics Reports **979**, 1 (2022). *Quantum magnonics: When magnon spintronics meets quantum information science*, H. Y. Yuan, Yunshan Cao, Akashdeep Kamra, Rembert A. Duine, and Peng Yan, Physics Reports **965**, 1 (2022).

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# **Coherent manipulation**

#### Normal mode splitting between a magnon (YIG sphere) and a cavity mode: Cavity-magnon polariton



They can drive cavity and magnon modes independently  $\rightarrow$  coherent superposition

Wolz et al, Communications Physics 3, 3 (2020)

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# **Coherent manipulation**



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# **Level repulsion and attraction**



$$E = \frac{E_1 + E_2}{2} \pm \sqrt{\left(\frac{E_1 - E_2}{2}\right)^2 + g^2}$$

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# **Dissipative coupling**

A magnetic sphere in a lossy Fabry-Perot cavity

#### Cavity photon spectrum:

- Localized waves at resonances
- Continuous travelling waves away from resonances

Coherent coupling to resonant modes – leads to level repulsion

Coupling to travelling modes – dissipative, leads to level attraction

- Non-Hermitian Hamiltonian
- Competition between level repulsion and attraction

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# **TUDelft**Level attraction

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M. Harder ... C.-M. Hu, Phys. Rev. Lett. 121, 137203 (2018) **18** Yaroslav M. Blanter ICTP Summer Scho



## Level attraction

#### A magnetic sphere in a lossy Fabry-Perot cavity



B. Yao, T. Yu, X. Zhang, W. Lu, Y. Gui, C.-M. Hu, and YMB, Phys. Rev. B **100,** 214426 (20) **19** Yaroslav M. Blanter ICTP Summer School, September 2023



# **Dissipative coupling**



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### Experiment: Excitation of spin waves

#### Magnetic nanowire array; FMR used to excite spin waves



#### Thickness: Not sufficient to support surface modes (20 nm)

J. Chen, T. Yu, C. Liu, T. Liu, M. Madami, Ka Shen, J. Zhang, S. Tu, M. Shah Alam, Ke Xia, M. Wu, G. Gubbiotti, YMB, G E. W. Bauer, and H. Yu, Phys. Rev. B **100,** 104427 (2019)

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J. Chen, T. Yu, C. Liu, T. Liu, M. Madami, Ka Shen, J. Zhang, S. Tu, M. Shah Alam, Ke Xia, M. Wu, G. Gubbiotti, YMB, G E. W. Bauer, and H. Yu, Phys. Rev. B **100**, 104427 (2019)

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### **Theoretical model**



We take for calculations s = 20 nm, a = 200 nm, d = 100 nm

Modes in the array: Close to the ferromagnetic resonance

$$\omega_{FMR} = \gamma \mu_0 \sqrt{H(H + M_0)}$$

Saturation magnatization

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## Spin waves in films



Spin waves can travel in both directions Spin-momentum locking: stray fields are polarized in a certain direction

$$B(k) \propto \left(m_z - i rac{k}{|k|} m_y
ight)$$

T. Yu, C. Liu, H. Yu, YMB, and G E. W. Bauer, Phys. Rev. B 99.134424 (2019)

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### **Dipolar interactions**

Free energy

 $F = -\mu_0 \int Mh^{dip} dr$  $h_{\beta}^{dip} = \frac{1}{4\pi} \partial_{\beta} \int dr' \frac{\partial_{\alpha} M_{\alpha}^{arr}(r')}{|r-r'|}$ 



- Only spin waves travelling in one direction generate dipolar fields above the film and thus couple to the array
- If the spin wave in the film is circularly polarized, it can only be excited in one direction
- Ellipticity leads to reduction of chirality





### Couplings





#### Parallel

Anti-parallel

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Coherent coupling of a magnon to a qubit via a cavity



The qubit is resonant with the magnon; the cavity is detuned and only serves to provide controlled interaction

Y. Tabuchi, S. Ishino, A. Noguchi, T. Ishikawa, R. Yamazaki, K. Usami, and Y. Nakamura, Science **349**, 405 (2015)

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# **Resonant coupling**



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# **Dispersive coupling**



Non-linear interaction (cross-Kerr):

$$\hat{H}_{int} = \tilde{K}\hat{m}^{\dagger}\hat{m}\hat{c}^{\dagger}\hat{c}$$

$$\hat{H}_q = \hbar \omega_q \hat{c}^\dagger \hat{c} + K \hat{c}^\dagger \hat{c}^\dagger \hat{c} \hat{c}$$

$$\hat{H}_m = \hbar \omega_m \hat{m}^\dagger \hat{m}$$

Shift of the frequency of one system depending on the state of another system





# **Resonant coupling**

Coherent coupling of a magnon to a qubit via a cavity



Strong coupling between a qubit and a magnon

Y. Tabuchi, S. Ishino, A. Noguchi, T. Ishikawa, R. Yamazaki, K. Usami, and Y. Nakamura, Science **349**, 405 (2015)

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# **Dispersive coupling**



Additions of magnons one by one lead to discrete shifts of the qubit frequency

D. Lachance-Quirion, Y. Tabuchi, S. Ishino, A. Noguchi, T. Ishikawa, R. Yamazaki, and Y. Nakamura, Sci. Adv. 3, e1603150 (2017)

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# **Creation of quantum states**

#### Resonant coupling regime



Da Xu, Xu-Ke Gu, He-Kang Li, Yuan-Chao Weng, Yi-Pu Wang, Jie Li, H. Wang, Shi-Yao Zhu, and J. Q. You, Phys. Rev. Lett. **130**, 193603 (2023)

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# **Magnon-qubit direct coupling**

#### Magnon lifetime: short (100 ns); difficult to create non-trivial quantum states

#### Idea: Use natural non-linear interactions



M. Kounalakis, G. E. W. Bauer, and YMB, Phys. Rev. Lett. 109, 037205 (2022)

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# **Magnon-qubit coupling**



M. Kounalakis, G. E. W. Bauer, and YMB, Phys. Rev. Lett. 109, 037205 (2022)

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# **Generation of cat states**

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M. Kounalakis, G. E. W. Bauer, and YMB, Phys. Rev. Lett. **109**, 037205 (2022) **35** Yaroslav M. Blanter ICTP Summer School, September 2023



### Conclusions

Spin waves interact with external magnetic field:

- This interaction can be strong in microwave cavities
- They facilitate manipulation of spin waves
- This may result in unusual behavior such as level attraction or chiral propagation
- Quantum properties of magnons can be detected by a qubit
  - One can bring magnons to non-trivial quantum states